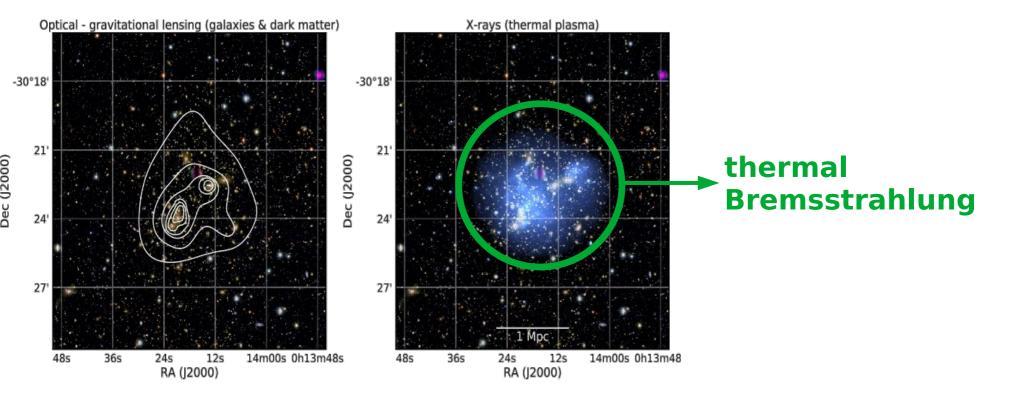
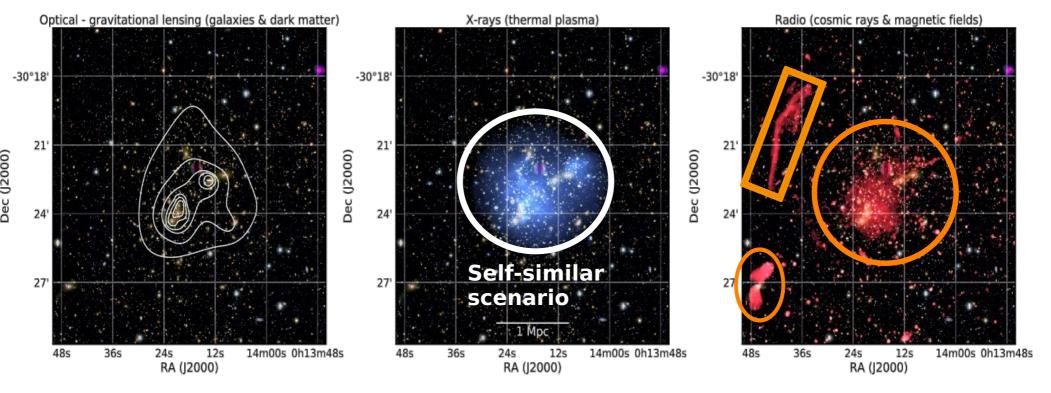


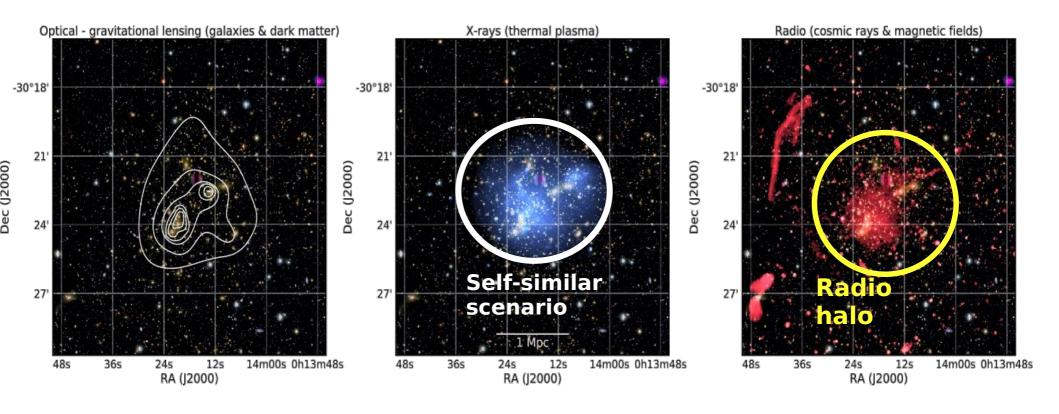
Clusters of galaxies are the largest objects in the Universe which are held together by their own gravity.

- Masses: $10^{14} 10^{15} M_{sun} (\sim 80\% \text{ dark matter})$
- Galaxies: 100s to 1000s
- Size: R~1 Mpc



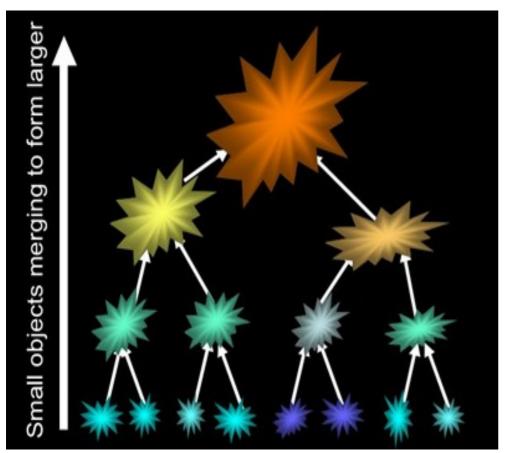
Synchrotron (CRe⁻ + weak B)

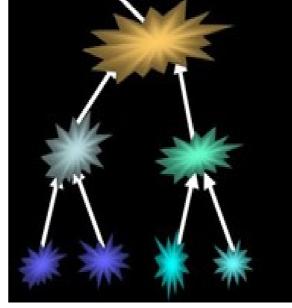




Self-similar scenario

• At cluster scales gravity dominates causing a self-similar evolution

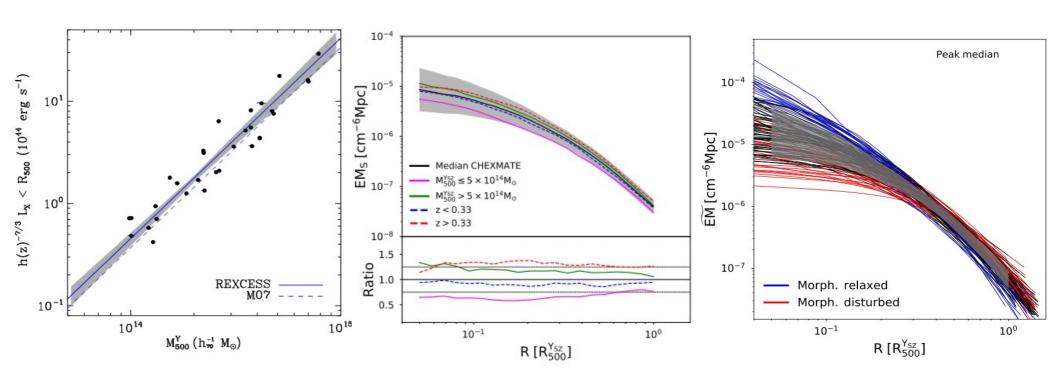






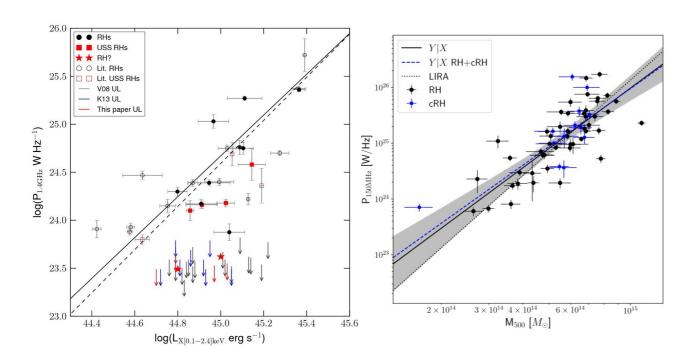
Self-similar scenario

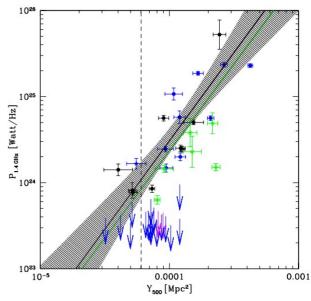
- At cluster scales gravity dominates causing a self-similar evolution
- Well studied through X-ray observations of ICM:
 - Scaling relations among global properties, L_x M, T M, L_x T (e.g. Pratt+09, Lovisari+2015) and resolved ones, $\rho(r)$, T(r) (e.g. Arnaud+2010, Rossetti+24)



Scaling laws in radio

• Non-thermal component scaling relations have been studied for integrated quantities: P_R





Scaling laws in radio

• Non-thermal component scaling relations have been studied for integrated quantities: P_R

• No systematic studies have been made on the scaling of spatially resolved properties (but also on the P_R redshift dependence)

Datasets



CHEX-MATE

LoTSS DR2

- Representative PSZ sample of 118 GC
- Homogeneous X-ray coverage
- Low and high redshift objects (Tier1 and Tier2)

Aims:

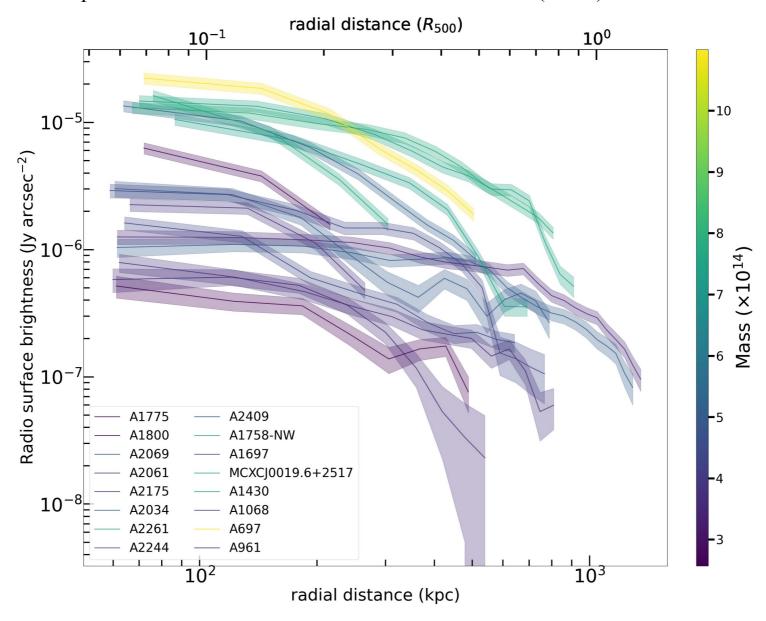
- Cluster absolute mass scale
- Cluster statistical properties
- How cluster properties changes over the time

- Deep 120-168 MHz survey of the Northen sky
- High sensitivity (100 muJy/beam)
- Ideal for halos studies (83 so far)

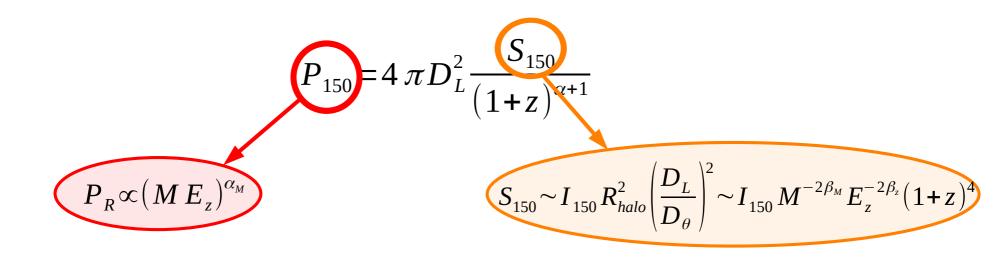
Botteon et al. 2022 Bruno et al. 2023 Zhang et al. 2023 Cassano et al. 2023 Cuciti et al. 2023 Jonese et al. 2023

CHEX-MATE Collaboration et al. 2021 Campitiello et al. 2022

• Extracted the radial profiles from the 16 CHEX-MATE - LoTSS DR2 (z<0.4)

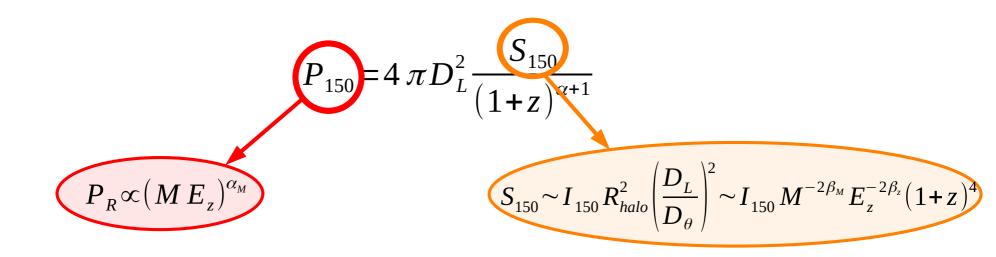


- Extracted the radial profiles from the 16 CHEX-MATE LoTSS DR2 (z<0.4)
- Derived the expected scaling in mass for the profiles assuming R_H^{\sim} M^{β} E_z^{β} . Compared the mass expected dependence with the best-fit scaling



$$I_{150}(r) \propto (M)^{\alpha_{M}-2\beta_{M}} E_{z}^{\alpha_{M}-2\beta_{z}} (1+z)^{-(3+\alpha)}$$

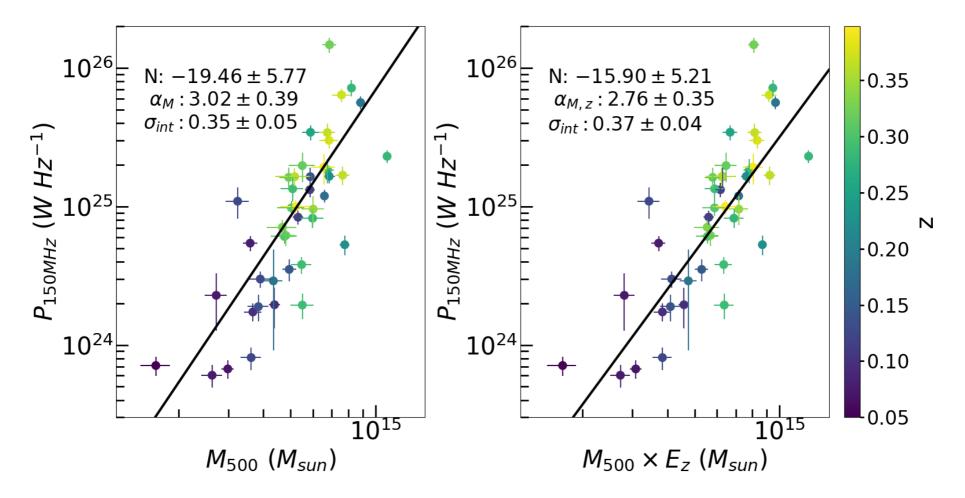
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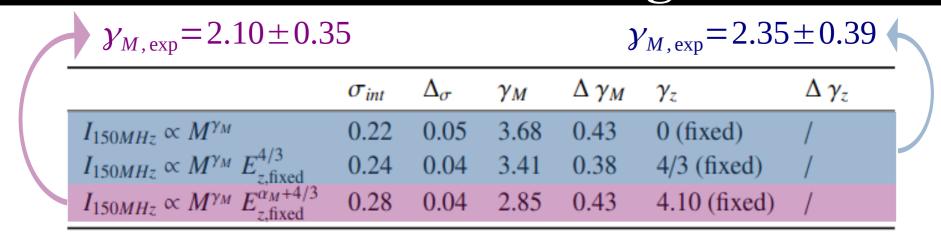


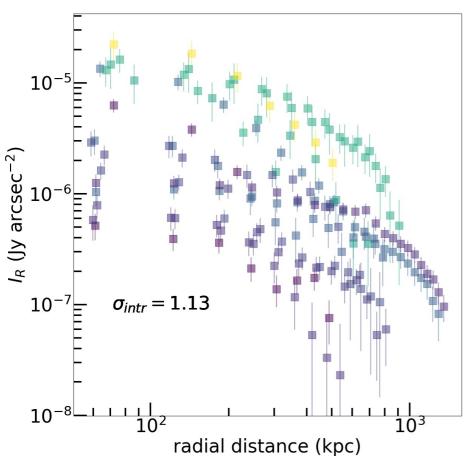
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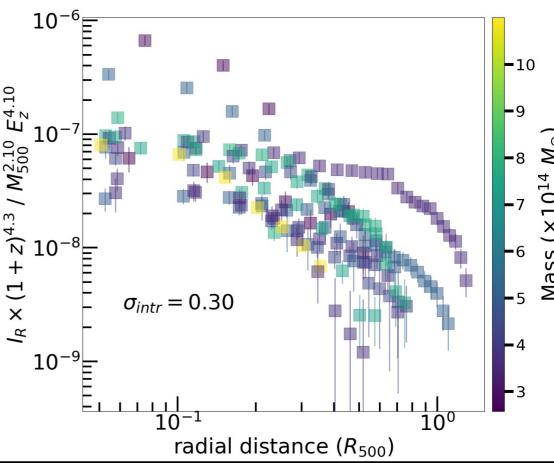
$$= I_{0} e^{\left(-A\frac{r}{r_{500}}\right)} M^{\gamma_{M}} E_{z}^{\gamma_{z}} (1+z)^{-(3+\alpha)}$$

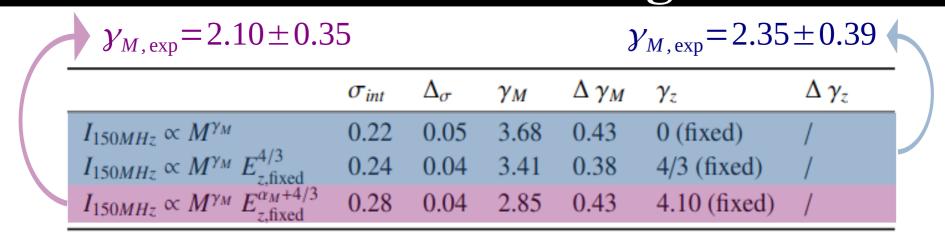
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- Exploited the Cuciti+23 RH sample to derive mass and redshift dependence

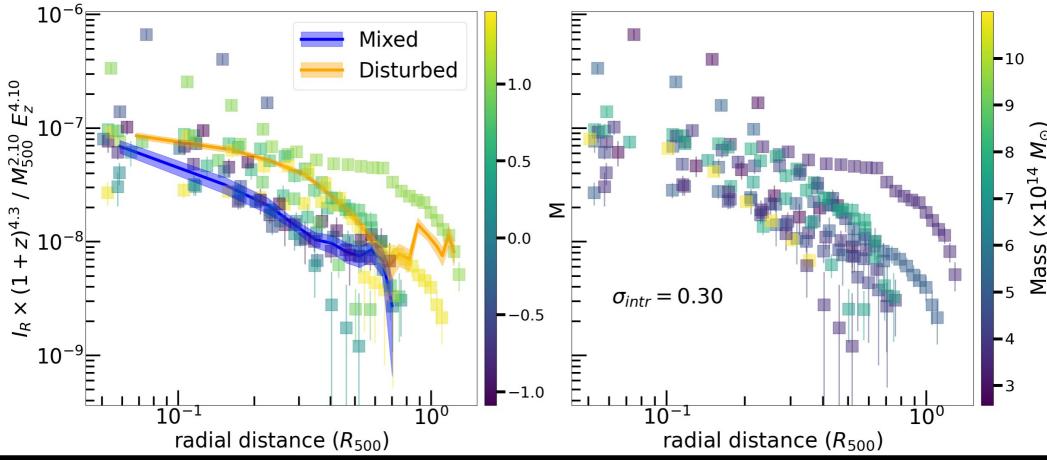










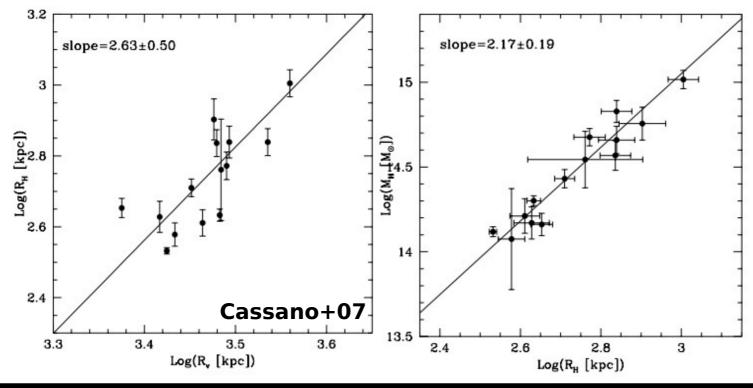


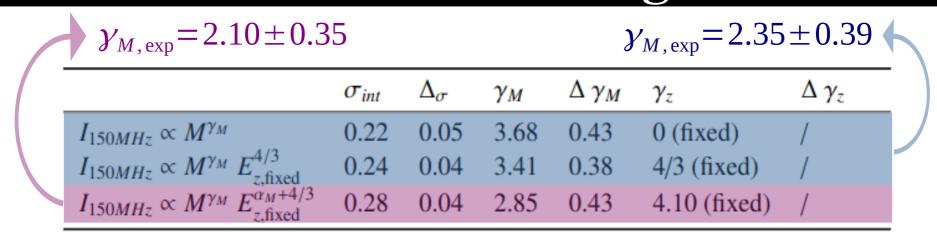
$$\gamma_{M, \, \text{exp}} = 2.10 \pm 0.35$$
 $\gamma_{M, \, \text{exp}} = 2.35 \pm 0.39$

$$\frac{\sigma_{int}}{\sigma_{int}} \Delta_{\sigma} \gamma_{M} \Delta_{\gamma_{M}} \gamma_{z} \Delta_{\gamma_{z}}$$

$$\frac{I_{150MHz} \propto M^{\gamma_{M}}}{I_{150MHz} \propto M^{\gamma_{M}}} \frac{0.22}{z_{\text{,fixed}}} \frac{0.05}{0.24} \frac{3.68}{0.04} \frac{0.43}{0.38} \frac{0 \, (\text{fixed})}{4/3 \, (\text{fixed})} / \frac{1_{150MHz} \propto M^{\gamma_{M}}}{z_{\text{,fixed}}} \frac{E_{z, \text{fixed}}^{\alpha_{M} + 4/3}}{0.28} \frac{0.04}{0.04} \frac{2.85}{0.43} \frac{0.43}{4.10 \, (\text{fixed})} / \frac{1_{150MHz}}{z_{\text{,fixed}}} \frac{M^{\gamma_{M}}}{z_{\text{,fixed}}} \frac{E_{z, \text{fixed}}^{\alpha_{M} + 4/3}}{0.28} \frac{0.04}{0.04} \frac{2.85}{0.43} \frac{0.43}{4.10 \, (\text{fixed})} / \frac{1_{150MHz}}{z_{\text{,fixed}}} \frac{M^{\gamma_{M}}}{z_{\text{,fixed}}} \frac{E_{z, \text{fixed}}^{\alpha_{M} + 4/3}}{0.28} \frac{0.04}{0.04} \frac{2.85}{0.43} \frac{0.43}{4.10 \, (\text{fixed})} / \frac{1_{150MHz}}{z_{\text{,fixed}}} \frac{M^{\gamma_{M}}}{z_{\text{,fixed}}} \frac{E_{z, \text{fixed}}}{0.28} \frac{0.04}{0.04} \frac{2.85}{0.43} \frac{0.43}{4.10 \, (\text{fixed})} / \frac{1_{150MHz}}{z_{\text{,fixed}}} \frac{1_{150MHz}}{z_{\text{,fixed$$

$$R_H \propto M^{\beta_M} E_z^{\beta_z}$$
 $\beta_M = 1/3, \beta_z = 0, -2/3$

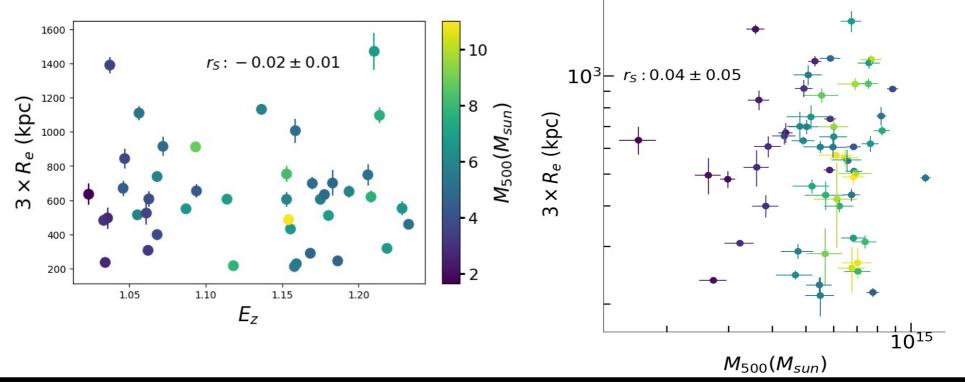


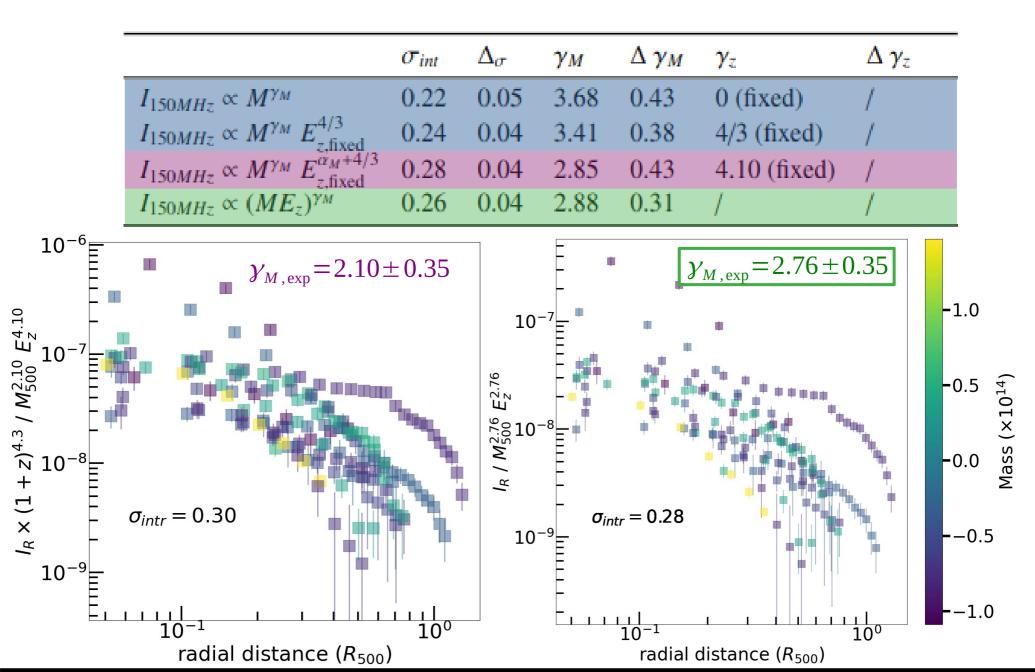


$$R_{\scriptscriptstyle H}\!\propto\! M^{eta_{\scriptscriptstyle M}} E_{\scriptscriptstyle z}^{eta_{\scriptscriptstyle z}}$$

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The radio halo power

$$P_R = f(M, z)$$

In the context of turbulent re-acceleration (Cassano+05,06)

$$P_R \propto M^{2-\Gamma} \frac{B^2 n_{e,r}}{(B^2 + B_{IC}^2)^2}$$

If
$$B = B_{M'} \left(\frac{M}{M'} \right)^b$$
:

$$\log(P_{150}) = N + \left(\frac{4}{3} - 2b\right) \log\left(\frac{M}{M'}\right) - 2\log\left[1 + \left(\frac{B_{CMB,0}}{B_{M'}}\right)^2 \frac{(1+z)^4}{(M/M')^{2b}}\right]$$

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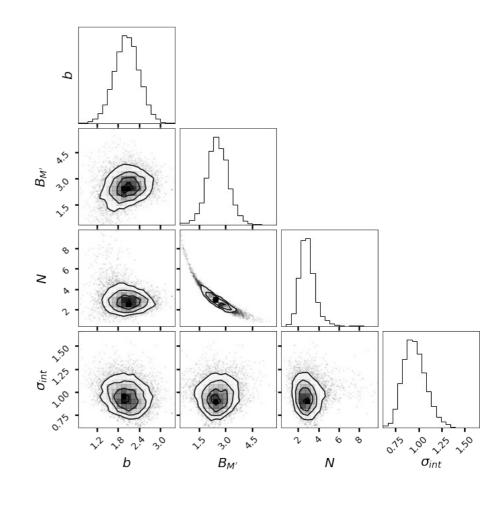
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Parameter	Prior	Best fit
b	U(0, 10)	2.05 ± 0.36
$B_{M'}$	U(0, 10)	2.53 ± 0.61
N	U(0, 10)	2.82 ± 0.83
σ_{int}	HC(0.5)	0.95 ± 0.13
b (fixed)	/	1.33
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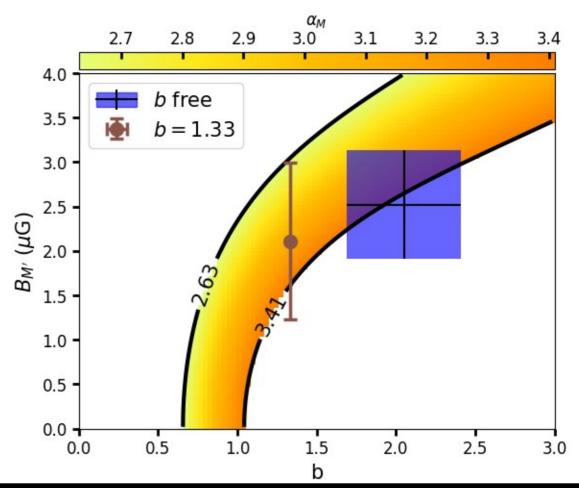


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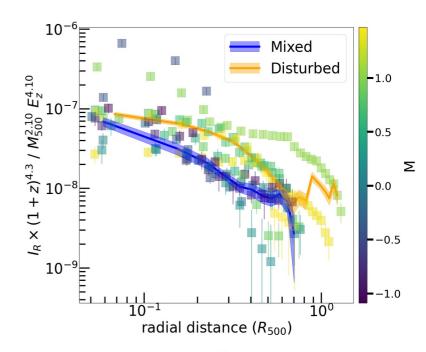
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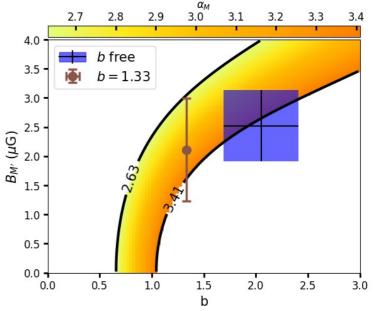


Summary

We studied the scaling properties of radio halo profiles finding:

- Strong cluster's mass depended of the profiles from I_R -M(z) relation
- The profile scatter to be singificantly reduced by M & z rescaling
- A residual dependence of the scaled profiles on the cluster dynamical state
- Evidence for no R_H-M relation, which, when taken into account allows to reconcile predictions and observations
- A strong (~2) mass dependence of the global cluster B



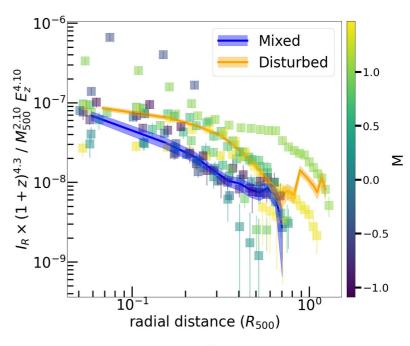


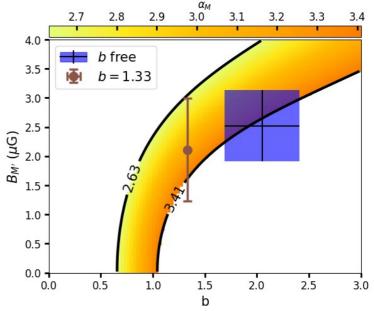
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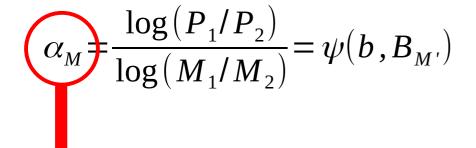
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Thank you for the attention!

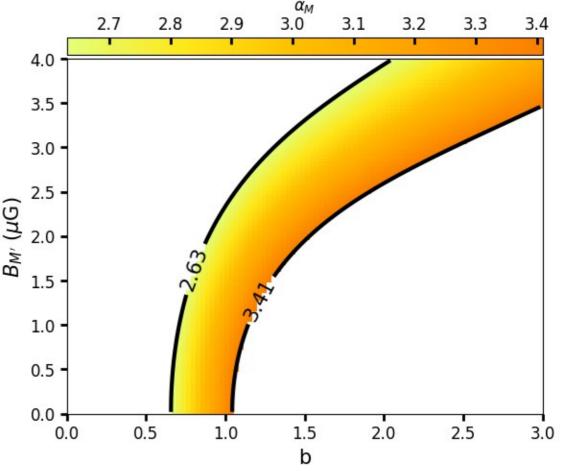




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Constrained from the observed P-M relation

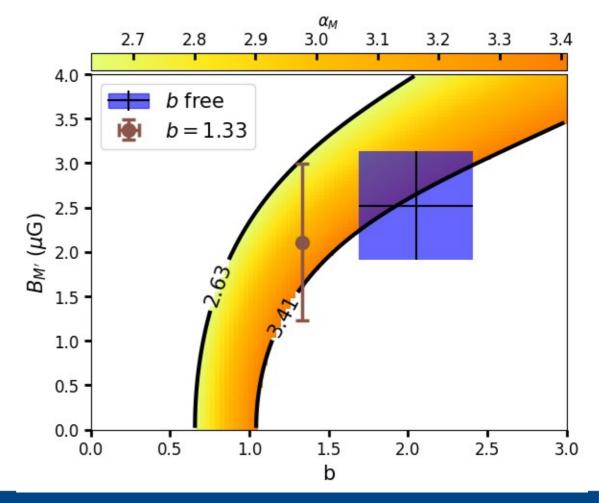


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